

## Photon Interaction, Buildup Factors and Fast Neutron Removal of Some Scintillators Used X-Ray and CT Imaging

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### Abstract

Photon and neutron shielding properties for scintillators used X-ray and CT imaging to investigate in this paper. Photon shielding properties of scintillators were investigated by computation for mass attenuation coefficient ( $\mu_m$ ), effective atomic number ( $Z_{eff}$ ), effective electron density ( $N_{el}$ ) and mean free path ( $MFP$ ) in the photon energy range 1 keV – 100 GeV, and buildup factor (BFs) were computed in photon energy region 15 keV – 15 MeV up to penetration depth of 40  $MFP$  by Geometric progression (G-P) fitting formula. Lastly, removal cross-section for fast neutron ( $\Sigma_R$ ) of scintillators were computed. The obtained results of some scintillators were compared, in terms of mass attenuation coefficient and mean free path with standard radiation shielding concretes. The results showed that the select scintillators better radiation shielding properties than the standard shielding concrete that can be developed as photon and neutron shielding materials.

**Keywords:** Scintillator; Radiation Shielding; Fast Neutron

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### Introduction

Today, radiation isotopes were used in many sectors such as agriculture and medicine whereas these sources are very dangerous which is including human life and sensitive laboratory device [1]. In the part of physics radiation, the attenuation of intensity radiation depends on the probability of interaction with the medium. These interactions are very important and can be discussed in the mass attenuation coefficient ( $\mu_m$ ) value.  $\mu_m$  value is basic quantity was used determined and explained the other values about the penetration of radiation in a medium such as the effective atomic number ( $Z_{eff}$ ), effective electron density ( $N_{el}$ ), mean free path ( $MFP$ ) and

buildup factors (*BFs*) [2, 3]. The probability of interaction with medium will higher when elements of the medium had a large atomic number and density, so, the choosing medium used to shield radiation should be selected medium as high atomic number and density elements [4 – 6].

Scintillator, one of material which absorbs radiation, high atomic number and density [7] and used in several applications such as medical imaging, the technology of astrophysical and space, nuclear fusion devices [8 – 11]. The utilization scintillator for radiation shielding and neutron is important which could be considered the basic value to discuss the penetration of radiation in a medium such as mass attenuation coefficient ( $\mu_m$ ), effective atomic number ( $Z_{eff}$ ), effective electron density ( $N_{el}$ ), half value layer (*HVL*) and build up factors (*BFs*) as mentioned before [12]. Gadolinium oxide ( $Gd_2O_3$ ) has properties such as high thermal and chemical stability, low phonon energy, non-toxicity, high refractive index, non-hygroscopic nature, high dielectric constant and high bandgap [13] and gadolinium oxysulfide ( $Gd_2OS$ ) has properties such as hard radiation stability, high X-ray absorption and high density makes it an effective trap of the incident X-ray photon [14] while gadolinium orthosilicate ( $Gd_2SiO_5$ : GSO) has high density, relatively fast decay time and high light output [15]. These properties of scintillators were according the properties of radiation shielding materials.

In this context was computed  $\mu_m$ ,  $Z_{eff}$ ,  $N_{el}$  and *MFP* value at photon energy ranging 1 keV – 100 GeV by WinXCOM software program while energy absorption buildup factors (*EABF*) and exposure buildup factors (*EBF*) of scintillator was computed at photon energy ranging 15 keV – 15 MeV by using G-P fitting method. Finally, to evaluate effective removal cross-sections values ( $\Sigma_R$ ) which presented neutron protection efficiency of scintillators.

## Theories and Computation Method

### *Scintillator*

Scintillator is inorganic material and used in field of ionizing nuclear radiation detection. To develop scintillator materials for radiation shield, the knowledge of radiation shielding parameters are very important for developed scintillators using against radiation such as  $\mu_m$  value due to this value show probability of interaction. In addition, when radiation interacts with medium,  $Z_{eff}$ ,  $N_{el}$ , *MFP*, *EABF* and *EBF* are basic values required to discuss the penetration of radiation in medium [12].

### *Radiation Shielding Properties*

The mass attenuation coefficient ( $\mu_m$ ) of scintillators was calculated theoretically by using mixture rule and WinXCOM software program [16, 17] as shown in equation (1);

$$\mu_m = \sum_{i=1}^n w_i \mu_{mi} \quad (1)$$

At this point,  $w_i$  and  $\mu_{mi}$  are weight fraction and mass attenuation coefficients of constituent elements, respectively. Using  $\mu_m$  value from equation (1) calculated  $Z_{eff}$  and  $N_{el}$  [18 – 20] as shown in equation (2) and (3);

$$Z_{eff} = \frac{\sigma_{t,a}}{\sigma_{t,el}} \quad (2)$$

$$N_{el} = \frac{Z_{eff} N_A}{M} \sum_i n_i \quad (3)$$

Effective atomic cross-section ( $\sigma_{t,a}$ ) were computed by equation (4);

$$\sigma_{t,a} = \frac{\sigma_t}{\sum_i n_i} \quad (4)$$

The total interaction cross-section ( $\sigma_t$ ) of scintillators have been computed with helped of  $\mu_m$  and estimated by equation (5);

$$\sigma_t = \frac{M \mu_m}{N_A} a \quad (5)$$

Total electron cross-section ( $\sigma_{t,el}$ ) were determined by equation (6);

$$\sigma_{t,el} = \frac{1}{N_A} \sum_i \frac{f_i A_i}{Z_i} (\mu_m)_i \quad (6)$$

At this point  $\sigma_{t,a}$ ,  $\sigma_t$  and  $\sigma_{t,el}$  are effective atomic cross-section, total photon interaction cross-section, and electronic cross-section, respectively.  $M$  is molecular weight,  $N_A$  is Avogadro's number,  $n_i$  is the number of formula units of the molecule,  $A_i$  indicates atomic weight of element  $i$ ,  $f_i$  presents a fractional abundance of constituent element and  $Z_i$  is an atomic number of element  $i$ . Mean free path ( $MFP$ ) of scintillators was evaluated by using linear attenuation coefficient ( $\mu$ ) [18, 19] as shown in equation (7);

$$MFP = \frac{1}{\mu} \quad (7)$$

#### *Build-Up Factors (BFs)*

The BFs are basically value used design medium for radiation shielding. BFs separated two types, 1) EABF and 2) EBF obtained by computing from Geometrical Progression (G-P) fitting method in energy region 15keV–15 MeV and these values can be following equation (8) – (9). Firstly, it is very important to know that the equivalent atomic number ( $Z_{eq}$ ) values must lie at specific energy between  $Z_1$  and  $Z_2$  atomic numbers ( $Z_1 < Z_{eq} < Z_2$ ) [21, 22] as shown in equation (8) – (9);

$$B(E, x) = 1 + \frac{b-1}{K-1} (K^x - 1), K \neq 1 \tag{8}$$

$$B(E, x) = 1 + (b-1)x, K = 1a \tag{9}$$

Herewith  $K(E, x)$ , photon dose multiplication factor,  $b$ , build-up factor corresponding to 1 *MFP* which obtained from equation (10);

$$K(E, x) = cx^a + d \frac{\tanh\left(\frac{x}{x_k} - 2\right) - \tanh(-2)}{1 - \tanh(-2)}, x \leq 40MFP \tag{10}$$

Lastly, fast neutrons removal cross-sections ( $\Sigma_R$ ) for scintillators [19, 22] can be computed using the following equation (11) and (12);

$$\Sigma_{\frac{R}{\rho}} = \sum_i w_i \left( \Sigma_{\frac{R}{\rho}} \right) \tag{11}$$

$$\Sigma_R = \sum_i \rho_i \left( \Sigma_{\frac{R}{\rho}} \right)_i \tag{12}$$

Where  $\rho_i$  is partial density and  $\Sigma_{R/\rho}$  are partial density and mass removal cross-section of *ith* element, respectively.

## Results and Discussion

The samples and properties of some scintillators are listed in Table 1. The  $\mu_m, Z_{eff}, N_{el}$  and *MFP* versus energy for some scintillators were presented in Fig. 1, 2, 3, and 4, respectively. The EABF and EBF values of some scintillators versus energy are presented in Fig. 6 and 7 (a) – (c) for different deep penetration at 1, 5, 10, 15, 20, 30, and 40 *MFP*. The EABF and EBF values for some scintillators versus deep penetration at energy 0.015, 0.15, 1.50, 3, 8, and 15 MeV are presented in Fig. 8 and 9 (a) – (d). Lastly, fast neutrons removal cross-sections ( $\Sigma_R$ ) for scintillators are shown in Fig. 9.

**Table 1** Properties of some scintillators used X-ray and CT imaging.

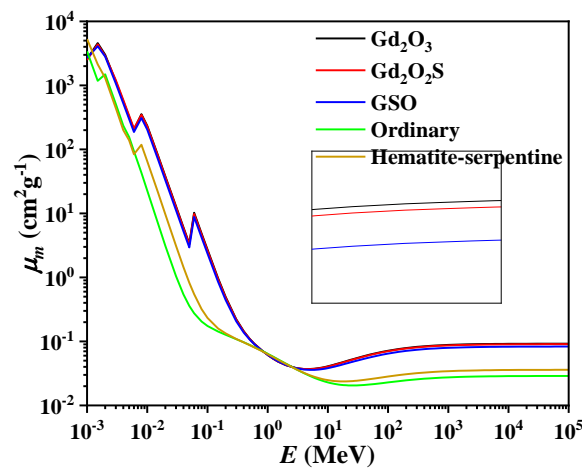
Sample	Density (g cc <sup>-1</sup> )	Ref.
Gd <sub>2</sub> O <sub>3</sub> :Eu <sup>3+</sup>	7.55	[11]
Gd <sub>2</sub> O <sub>2</sub> S:Pe,Ce,F	7.34	[11]
GSO:Ce	6.70	[11]

$\mu_m$ ,  $Z_{eff}$ ,  $N_{el}$  and MFP of Scintillators

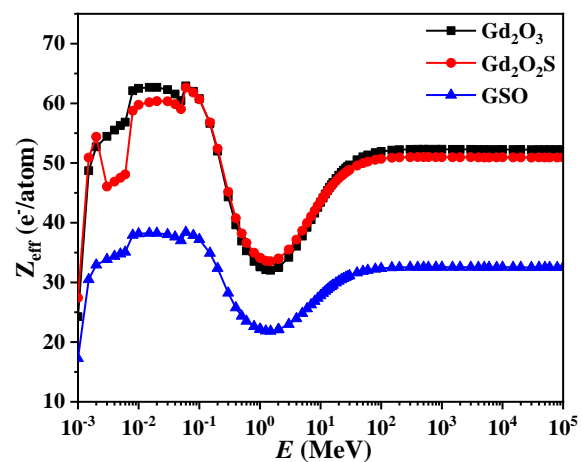
The  $\mu_m$  versus photon energy at 1 keV–100 GeV of scintillators is presented in Fig. 1. From Fig. 1  $\mu_m$  values of all scintillators are decrease exponentially with increasing photon energies. At low energy,  $\mu_m$  values are very high and decrease quickly with increasing energies. At intermediate energy,  $\mu_m$  values are slow rate and becoming approximately constant at high energy. These events were discussed on the photoelectric effect, Compton scattering and pair production which appear and dominate at low, intermediate and high energy, respectively. From Fig. 1 graph are not continuous because of absorption edges of elements as presented in Table 2 and  $Gd_2O_3$  scintillator has higher  $\mu_m$  than other scintillators as indicated that  $Gd_2O_3$  was used radiation shielding better other scintillators.

**Table 2** Absorption edges (keV) of elements.

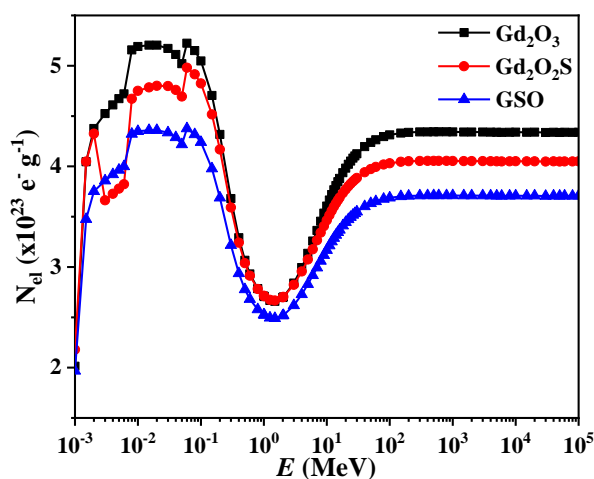
Element	K	L1	L2	L3	M1	M2	M3	M4	M5
Si	1.839	-	-	-	-	-	-	-	-
S	2.472	-	-	-	-	-	-	-	-
Gd	50.24	8.38	7.93	7.24	1.88	1.69	1.54	1.22	1.19



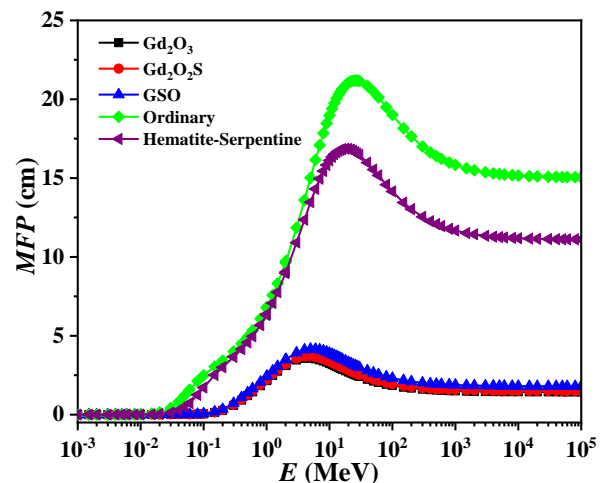
**Fig. 1**  $\mu_m$  versus photon energy.



**Fig. 2**  $Z_{eff}$  versus photon energy.



**Fig. 3**  $N_{el}$  versus photon energy.



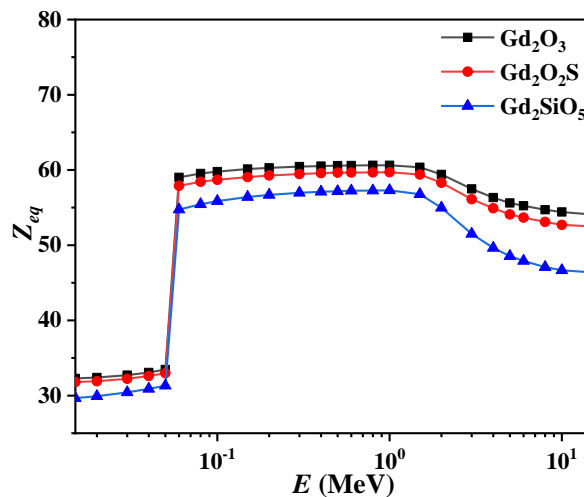
**Fig. 4** MFP versus photon energy.

The  $Z_{eff}$  versus photon energy for scintillators has been presented in Fig. 2. It has been observed in  $Z_{eff}$  because of the main of different photon interaction processes. At lower energies, the photoelectric effect is the main photon interaction. When energy is adjacent to K-absorption edges of each element (Si, S, Gd) existent in scintillators,  $Z_{eff}$  values sudden shift like  $\mu_m$  which indicates that correspond to elements K-absorption edge. This mechanism is main below 0.10 MeV. Hence, this event is computed by  $Z$  dependence of total atomic cross-sections and high  $Z$  elements from cross-section of the photoelectric absorption process. At intermediate energy ranging, the main interaction is Compton scattering,  $Z_{eff}$  value of scintillators was found the lowest. And then at high energy ranging which pair production is main interaction,  $Z_{eff}$  value increases again. At energy more than over about 100 MeV,  $Z_{eff}$  value nearly constant, this is discussing from cross-section of which possesses a weaker  $Z^2$  dependence. However, cross-section of Compton scattering depends on  $Z$  whereas less than the photoelectric absorption and pair production process. The  $N_{el}$  value has the same trend of  $Z_{eff}$  that indicates the behavior of  $N_{el}$  like  $Z_{eff}$  as shown in Fig.3.

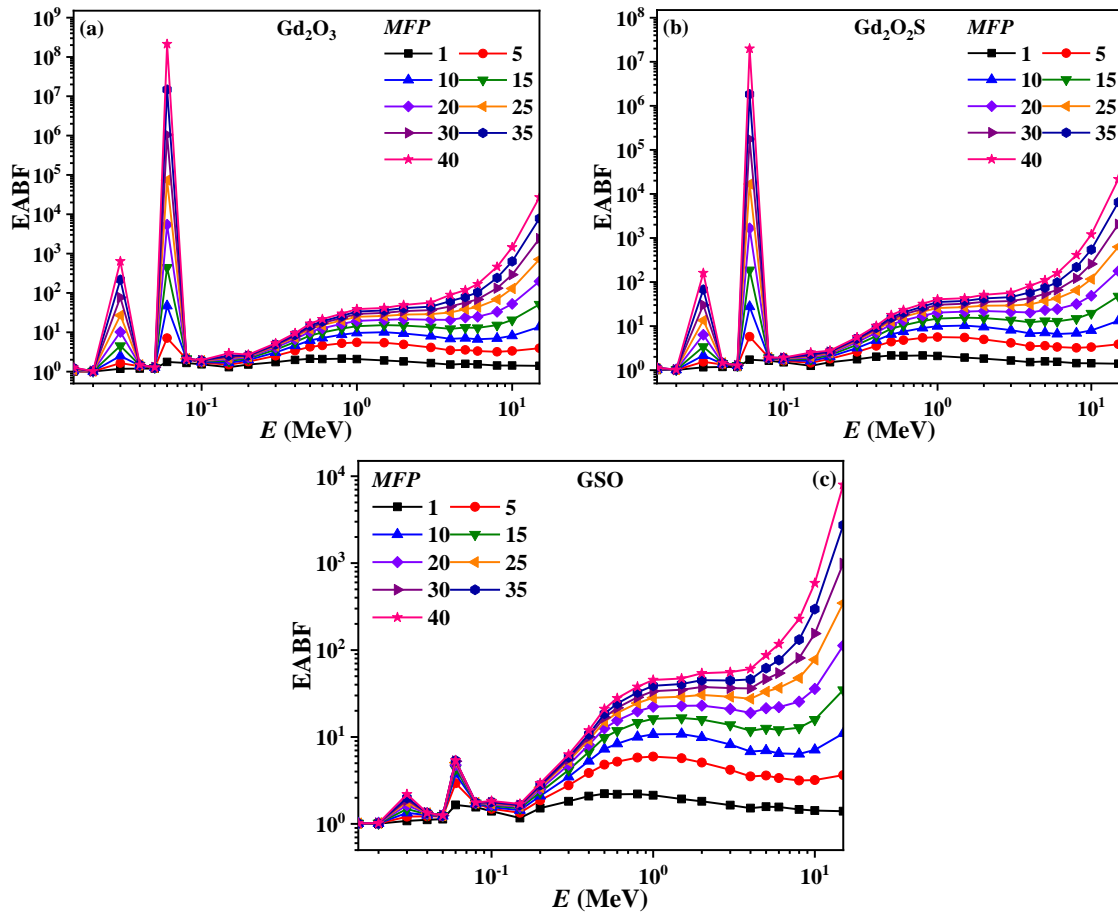
The  $MFP$  value for c scintillators are exhibited in Fig. 4. As indicated  $MFP$  value increase with increasing energy while GSO and  $Gd_2O_3$  have the highest and lowest  $MFP$ , respectively. This results from the density of scintillators which high density presented radiation absorption efficiency of scintillator.  $Gd_2O_3$  has the highest density, thus the density of scintillator plays an important role in affecting photon attenuation. Since  $Gd_2O_3$  has the lowest  $MFP$  at the same energy, so  $Gd_2O_3$  is required a small thickness when compared with the other scintillator and used for radiation shielding.

*EABF, EBF and  $\Sigma_R$  of Scintillators*

EABF and EBF of scintillators have been explained depending on density, photon energy, and deep penetration. Fig. 5 shows  $Z_{eq}$  of scintillators versus photon energy. It is clearly seen that  $Gd_2O_3$  scintillator has the highest  $Z_{eq}$  values.

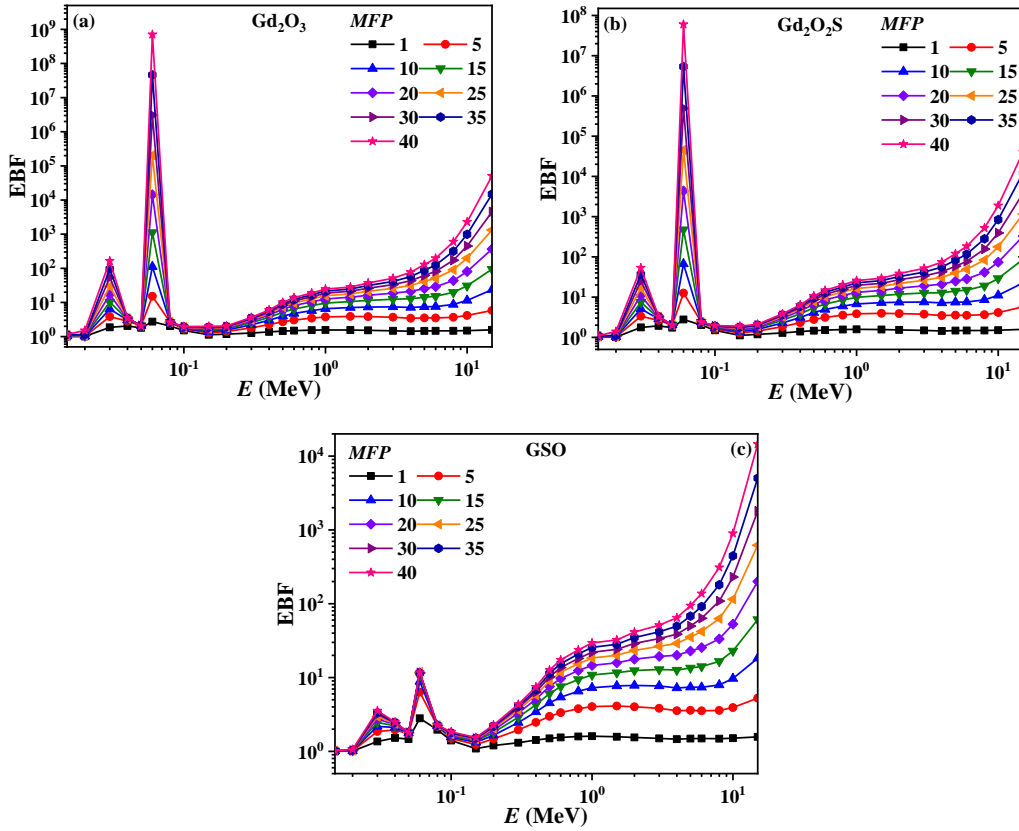


**Fig. 5**  $Z_{eq}$  versus photon energy.

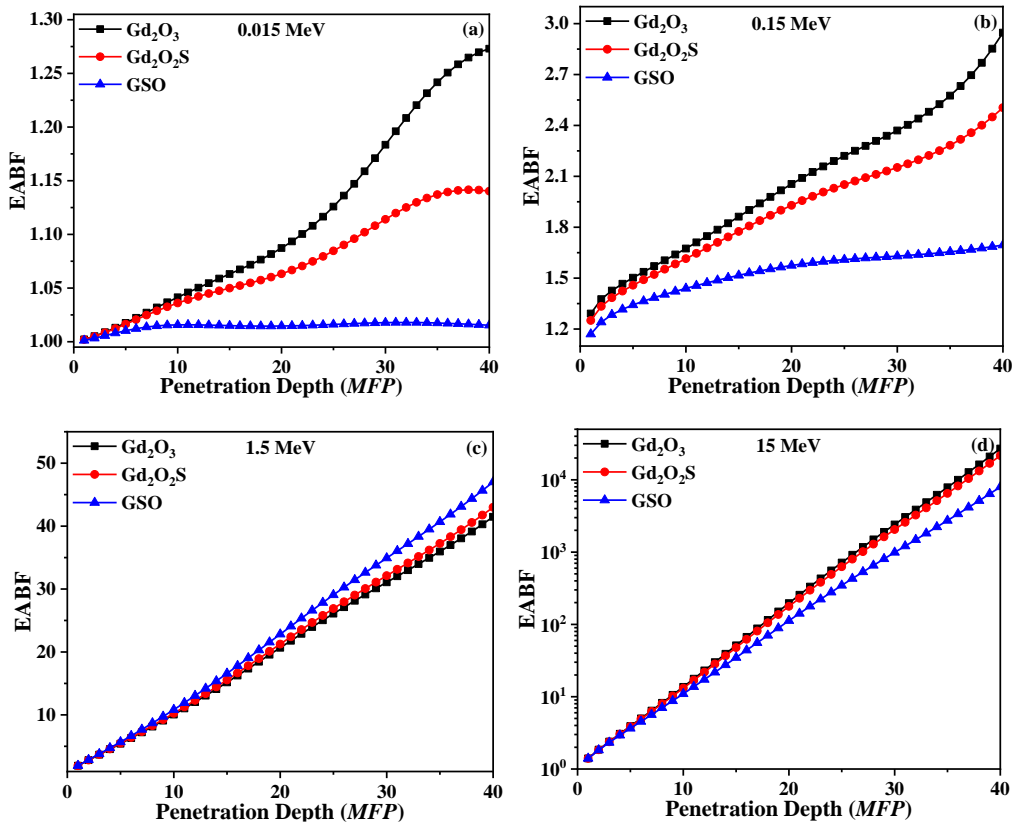


**Fig. 6** EABF versus photon energy at 1, 5, 10, 15, 20, 25, 30, 35, and 40 *MFP* of (a) Gd<sub>2</sub>O<sub>3</sub>, (b) Gd<sub>2</sub>O<sub>2</sub>S, and (c) GSO.

The EABF and EBF, with energy, for all scintillators at different deep penetration, have been exhibited in Fig. 6 and 7(a) – (c), respectively. These figures, EABF and EBF values for all scintillators increase with increasing photon energy and the maximum values indicated that EABF and EBF are dependent on deep penetration and composition of scintillators. These events can be discussed on the important interaction. At low energy ranging, values of EABF and EBF are smallest because of photons were completely absorbed and the main interaction process is the photoelectric effect. At intermediate energy ranging, the values of EABF and EBF are highest because of Compton scattering is the main interaction process. At high energy ranging, pair production is main interaction, photons have been absorbed again. The EABF and EBF results of scintillators have two sharp peaks at 40 and 60 keV which may be due to K-absorption edges of elements. The highest EABF and EBF values were found at 40 *MFP* deep penetration while the lowest values were found at 1 *MFP* this because of multiple scatterings were occurred at high deep penetration.

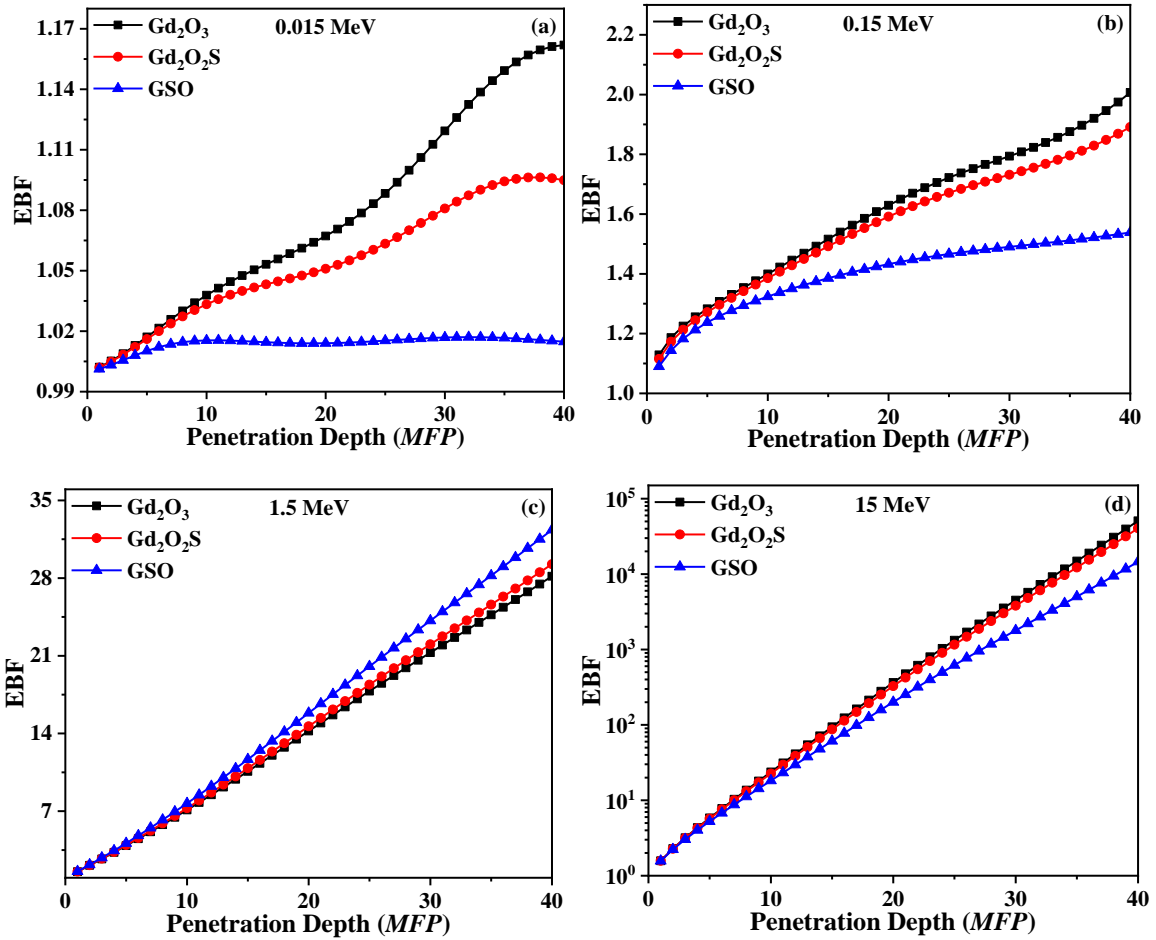


**Fig. 7** EBF versus photon energy at 1, 5, 10, 15, 20, 25, 30, 35, and 40 MFP of (a)  $Gd_2O_3$ , (b)  $Gd_2O_2S$ , and (c) GSO.



**Fig. 8** EABF versus deep penetration for (a) 0.015 MeV, (b) 0.15MeV, (c) 1.50 MeV and (d) 15 MeV.

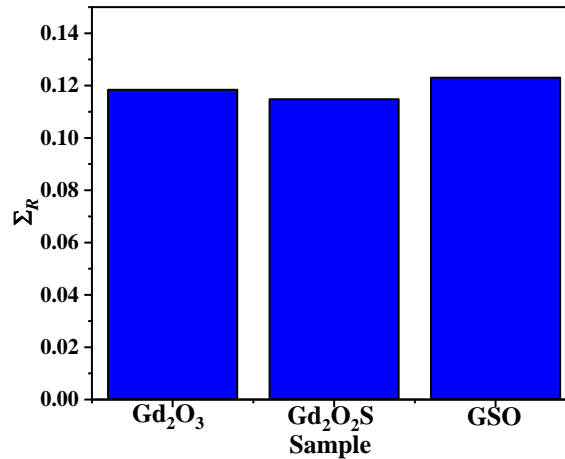




**Fig. 9** EBF versus deep penetration for (a) 0.015 MeV, (b) 0.15 MeV, (c) 1.50 MeV, and (d) 15 MeV.

The EABF and EBF, with deep penetration, for all scintillators at energy 0.015, 0.15, 1.50, and 15 MeV have been exhibited in Fig 8, and 9 (a) – (d), respectively. It can be noted that EABF and EBF value of scintillators increase with increasing energy and deep penetration. These events can be discussed on the basis of the partial photon interaction process. At low, intermediate and high energy ranging, photoelectric effect, Compton scattering, and pair production is the main interaction, respectively. At energy 0.015 and 0.15 MeV, EABF and EBF of Gd<sub>2</sub>O<sub>3</sub> value presented maximum values. This indicates that at low energies, EABF and EBF depend clearly on the chemical composition of scintillators. At 15 MeV energy, Gd<sub>2</sub>O<sub>3</sub> showed the largest EABF and EBF again that because of pair/triplet production in scintillators which generates electron/positron pair. This process becomes to allow photon to accumulate in Gd<sub>2</sub>O<sub>3</sub>.

The removal cross-section for fast neutron ( $\Sigma_R$ ) of scintillators is presented in Fig. 8. The highest  $\Sigma_R$  (cm<sup>-1</sup>) value was found for GSO, this reveals that low-Z elements are not only responsible for neutron removal but also including of low and high-Z elements [23].



**Fig. 10**  $\Sigma_R$  of scintillators.

## Conclusion

In this context, scintillators used X-ray and CT imaging were estimated radiation shielding at photon energy ranging 1 keV – 100 GeV, buildup factors at photon energy range 15 keV – 15 MeV up to deep penetration of 40 *MFP*, and fast neutron removal cross-section. The simulated values find that Gd<sub>2</sub>O<sub>3</sub> scintillator showed excellent radiation shielding because of the highest values for mass attenuation coefficient ( $\mu_m$ ), effective atomic number ( $Z_{eff}$ ) and effective electron density ( $N_{el}$ ) and lower values for mean free path (*MFP*) and buildup factors (*BFs*). While neutrons shielding, GSO scintillator showed excellent because of the highest removal cross-section for fast neutron ( $\Sigma_R$ ) value. This study on the interaction of radiation and neutron with scintillators provides sufficient information for the selection of these scintillators for the design and applications of shielding materials in various applications.

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